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## Dual solutions of cross diffusion effects on MHD Peristaltic flow in a conduit

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Abstract. Thepurpose of this study is to analyze the influence of wall flexibility and cross diffusion on the peristaltic transport of a MHDdissipative fluid in uniform channel in the presence of Joule heating. Dual solutions are obtained for aligned and transverse magnetic field cases. The exact solution is obtained for stream function and velocity. Further, shooting process isemployedtosolvetheenergy and concentration equations. The influence of main parameters on the present flow is explained graphically. The increase in magnetic field decreases the velocity of the fluid. Also Dufour and Soret numbers enhances the temperature while they reduses the concentration field.

### Nomenclature:

- Cartesian coordinates *x*, *v*
- fluid velocities u, v
- pressure p
- viscosity μ
- mean width of the channel d
- amplitude а
- wavelength λ
- wave speed С
- specific heat  $C_p$
- kinematic viscosity υ
- thermal conductivity  $k_0$
- concentration susceptibility  $C_{s}$
- electrical conductivity  $\sigma$
- magnetic field  $B_0$
- Т temperature
- amplitude ratio ε

- density of the fluid ρ
- time t
- concentration С
- $C_0, C_1$  concentrations fluid at lowerand upper walls

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- $K_T$ diffusion ratio
- $D_{R}$ mass diffusion coefficient
- Sr Soret parameter
- Du Dufour parameter
- yield stress  $au_0$
- $T_{0}, T_{1}$ temperatures of fluid at lower and upper walls
- $E_1, E_2, E_3$  elasticity parameters
- $\theta$ Non-dimensional temperature
- Ø Non-dimensional concentration
- magnetic parameter Μ
- Nu Nusselt number
- Brinkman number Br

Ec

Eckert number

 $\delta$  wave number

 $\psi$  stream function

#### **1. Introduction**

Peristalsis is an importantmechanismoffluid transport which is formed by the propagation of a peristaltic wave on the channel walls. We can observe this natural phenomenon in many physiological systems such as swallowing foodstuff, passage of urine, transport of lymph, blood pumping in blood vessels, ovum transport etc. Also it is useful in pumping some industrial harsh fluids. Further engineers are adapting this principle in preparing different models in pumping michanary. In view of these real applications many authors have investigated the peristaltic transport in various circumstances [1-4].

The influence of magnetic field,heat and mass transfer on peristalsis has huge applications in industrial and biomedical endineering like drug distribution schemes, canser treatment, reduction of blood loss during operations, design of heat exchengers, solar energy and many others. Few authors have studied the combined effects of heat transfer and magnetohydrodynamics on peristaltic flow. Some of the recent studies are presented in references [5-16]. Very recently, the researchers [17,18] analyzed the convective heat transfer in MHD flows by considering the cross diffusion effects.

Motivated by the above discussed applications and studies, we examined the effects of Joul heating, Dufour and Soret numbers on peristaltic transport of a viscous fluid in an elastic channel in the presence of inclined magnetic field. The analytical solutions are presented for stream function and velocity and the energy and concentration equations are solved numerically. The impact of key parameters on the flow characteristics is discussed through graphs.

## 2. Modelling of the problem

We consider the peristaltic transport of a viscous fluid in a two dimensional uniformchannel. The flow is generated by a peristaltic wave propagating along the elastic walls of the channel with a constant speed c (Fig.1). The Joule heating, diffusionand dissipation effects are considered. The geometry of the channel is given as

$$\bar{h}(\bar{x},\bar{t}) = a \sin \frac{2\pi}{\lambda} (\bar{x} - c\bar{t}) + d$$
(1)

As per the assumptions made above, the simplified non-dimensional governing equations and corresponding boundary conditions for the present study are given by (See the references [9-13])



Figure 1.Flow configuration



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$$\frac{\partial^2 \theta}{\partial y^2} + Br \left[ \left( \frac{\partial^2 \psi}{\partial y^2} \right)^2 + M^2 \cos^2 \alpha \left( \frac{\partial \psi}{\partial y} \right)^2 \right] + \Pr Du \frac{\partial^2 \phi}{\partial y^2} = 0, \tag{4}$$

$$\frac{\partial^2 \phi}{\partial y^2} + ScSr \frac{\partial^2 \theta}{\partial y^2} = 0, \tag{5}$$

Appropriate boundary conditions are

$$\psi = 0, \frac{\partial^2 \psi}{\partial y^2} = \frac{\partial \theta}{\partial y} = \frac{\partial \phi}{\partial y} = 0, \text{ at } y = 0$$
 (6)

$$\frac{\partial \psi}{\partial y} = 0, \ \theta = 1, \ \phi = 1 \text{ at } y = h = [\varepsilon \sin 2\pi (x - t) + 1]$$
(7)

$$\frac{\partial^3 \psi}{\partial y^3} - M^2 \frac{\partial \psi}{\partial y} \cos^2 \alpha - \left( E_3 \frac{\partial^2 h}{\partial x \partial t} + E_1 \frac{\partial^3 h}{\partial x^3} + E_2 \frac{\partial^3 h}{\partial x \partial t^2} \right) = 0, \text{ at } y = h$$
(8)

The dimensionless parameters used in the above simplified equations are :

$$x\lambda = \bar{x}, yd = \bar{y}, \psi c d = \bar{\psi}, uc = \bar{u}, vc = \bar{v}, p = \frac{d^2 \bar{p}}{\mu c \lambda}, \tau_{xy} = \frac{d \bar{\tau}_{xy}}{\mu c}, t = \frac{c\bar{t}}{\lambda}, \delta = \frac{d}{\lambda}, \\ \varepsilon = \frac{a}{d}, \operatorname{Re} = \frac{\rho c d}{\mu}, \theta = \frac{(T - T_0)}{(T_1 - T_0)}, \phi = \frac{(C - C_0)}{(C_1 - C_0)}, Sc = \frac{\upsilon}{D_B}, Sr = \frac{\rho D_B K_T}{\mu (C_1 - C_0)}, \\ M = \sqrt{\frac{\sigma}{\mu}} B_0, \operatorname{Pr} = \frac{\rho \upsilon c_f}{k_0}, Ec = \frac{c^2}{c_f (T_1 - T_0)}, Du = \frac{D_B K_T (C_1 - C_0)}{\mu c_p c_s (T_1 - T_0)}, \\ E_1 = \frac{-\tau d^3}{\lambda^3 \mu c}, E_2 = \frac{m_1 c d^3}{\lambda^3 \mu}, E_3 = \frac{c d^3}{\lambda^2 \mu}, h = \frac{\bar{h}}{d} = 1 + \varepsilon \sin 2\pi (x - t), \end{cases}$$
(9)

#### **3.** Solution of the problem

By differentiating equation (2) with respect to y we obtain

$$\frac{\partial^4 \psi}{\partial y^4} - M^2 \frac{\partial^2 \psi}{\partial y^2} \cos^2 \alpha = 0 \tag{10}$$

By solving equation (10) with the boundary conditions (6), (7) and (8) we attain the stream function and the velocity as

$$\psi = \frac{C_1}{M^2 \cos^2 \alpha} \left[ \frac{\sinh(M \cos \alpha) y}{M \cos \alpha \cosh(M \cos \alpha) h} - y \right]$$
(11)

$$u = \frac{C_1}{M^2 \cos^2 \alpha} \left[ \frac{\cosh(M \cos \alpha) y}{\cosh(M \cos \alpha) h} - 1 \right]$$
(12)

Where

$$C_{1} = -8\varepsilon \pi^{3} \left[ (E_{1} + E_{2})\cos 2\pi (x - t) - \frac{E_{3}}{2\pi} \sin 2\pi (x - t) \right]$$

The equations (4) and (5) are coupled and nonlinear therefore these equations are solved numerically by employing R-K-Fehlberg integration schemewith the help of (7), (8) and (11). The Nusselt and Sherwood numbers at the wall are defined by

$$Nu = -\left(\frac{d\theta}{dy}\right)_{at \ y=h}, \ Sh = -\left(\frac{d\phi}{dy}\right)_{at \ y=h}$$
(13)



#### 4. Results of the problem

The impact of relevant parameters on the flow characteristics are

Presented with the help of graphical and tabular representations. In this section we used the fixed values of various parameters as x = 0.2,  $\varepsilon = 0.2$ , t = 0.1, M = 2,  $E_1 = 0.2$ ,  $E_2 = 0.2$ ,  $E_3 = 0.1$ , Pr = 5, Br = 0.2, Sc = 0.6, Sr = 0.2, Du = 0.2,  $\alpha = 0$  (Normal magnetic field) and  $\alpha = \pi/6$  (Inclined magnetic field).



Figures 2 and 3 are plotted to explain the effects of amplitude ratio  $\varepsilon$  and magnetic parameter M on velocity. We notice that velocity is an increasing function  $\varepsilon$  and decreasing function of M. It is clear that the velocity field of the fluid effected by the magnetic force. Figures 4 - 9 are shown that the wall tension  $(E_1)$  and mass characterization  $(E_2)$  parameters enhances the velocity and temperature while they reduces the concentration. Further, opposite behavior is identified for damping force parameter  $E_3$ .



**Figure 9.** Concentration profiles for  $E_3$ 



Figs. 10-13 illustrates the influence of Dufour and Soret parameters on thermal and concentration fields. It is clear that the rising values of cross diffusion parameters boosts the thermal field and supresses the concentration field. It is interesting to mention that the effect of cross diffusion is high when the magnetic field is transverse. It is noticed from the Figs. 15 and 16 that the increasing values of Brickman number encourages the thermal field and declines the concentration field in both cases.



Figure 13. Concentration profiles for Sr



Figure 15. Concentrationprofiles for Br

Table 1 Variations in physical quantities in uniform and non-uniform channel cases

					Aligned magnetic field		Normal magnetic field	
$E_{1}, E_{2}$	$E_3$	Du	Sr	Br	$-\theta'(h)$	$-\phi'(h)$	$-\theta'(h)$	$-\phi'(h)$
0.1					-0.163415	-0.166791	-0.199292	-0.203410
0.15					-0.173566	-0.176963	-0.211671	-0.215814
0.2					-0.183778	-0.187197	-0.224126	-0.228295
	0.1				-0.144535	-0.147522	-0.176267	-0.179910
	0.2				-0.153513	-0.156518	-0.187216	-0.190881
	0.3				-0.162546	-0.165570	-0.198233	-0.201920
		0.1			-0.186978	-0.191365	-0.228028	-0.233378
		0.4			-0.200235	-0.204717	-0.244195	-0.249661
		0.7			-0.213777	-0.218356	-0.260711	-0.266295
			0.4		-0.768848	-0.787342	-0.937645	-0.960199
			0.6		-0.824784	-0.843735	-1.005861	-1.028972
			0.8		-0.882102	-0.901521	-1.075762	-1.099445
				0.1	-1.634147	-1.667914	-1.992913	-2.034093
				0.15	-1.735654	-1.769628	-2.116705	-2.158137
				0.2	-1.837783	-1.871965	-2.241254	-2.282940

It is evident from the table 1 that the cross diffusion and Brickman number have tendency to reduce the heat and mass transfer rate. The similar trend has been observed for rising values of wall tension (  $E_1$ ), mass characterization ( $E_2$ ) and damping force ( $E_3$ ) parameters.

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